Time-dependent density-functional formalism
(first: electrons only, nuclei are fixed or treated classically)

Basic 1-1 correspondence:

$$v(rt) \xleftrightarrow{\;1\text{-}1\;} \rho(rt)$$

The time-dependent density determines uniquely the time-dependent external potential and hence all physical observables for fixed initial state.

KS theorem:

The time-dependent density of the **interacting** system of interest can be calculated as density

$$\rho(rt) = \sum_{j=1}^{N} \left| \varphi_j(rt) \right|^2$$

of an auxiliary **non-interacting** (KS) system

$$i\hbar \frac{\partial}{\partial t} \varphi_j(rt) = \left(-\frac{\hbar^2 \nabla^2}{2m} + v_S[\rho](rt) \right) \varphi_j(rt)$$

with the **local** potential

$$v_S[\rho(r't')](rt) = v(rt) + \int d^3r' \frac{\rho(r't)}{|r-r'|} + v_{xc}[\rho(r't')](rt)$$

The TDKS equations follow (like in the static case) from:

i. the basic 1-1 mapping and

ii. the TD V-representability theorem (R. van Leeuwen, PRL <u>82</u>, 3863 (1999)).

A TDDFT variational principle exists as well, but this is more tricky (R. van Leeuwen, PRL <u>80</u>, 1280 (1998)).

Proof of basic 1-1 correspondence between $v(\vec{r}t)$ **and** $\rho(\vec{r}t)$

define maps

$$F: \quad v(\vec{r}t) \mapsto \Psi(t)$$

$$\tilde{F}: \quad \Psi(t) \mapsto \rho(\vec{r}t)$$

G

potentials $v(\vec{r}t)$

F
solve tdSE
with <u>fixed</u>
$\Psi(t_o) = \Psi_o$

wave functions $\Psi(t)$

$\tilde{F}$

$\rho(\vec{r}t) = \langle \Psi(t)|\hat{\rho}(\vec{r})|\Psi(t)\rangle$

$\hat{\rho}(\vec{r}) = \sum_s \hat{\psi}_s^+(\vec{r})\hat{\psi}_s(\vec{r})$

densities $\rho(\vec{r}t)$
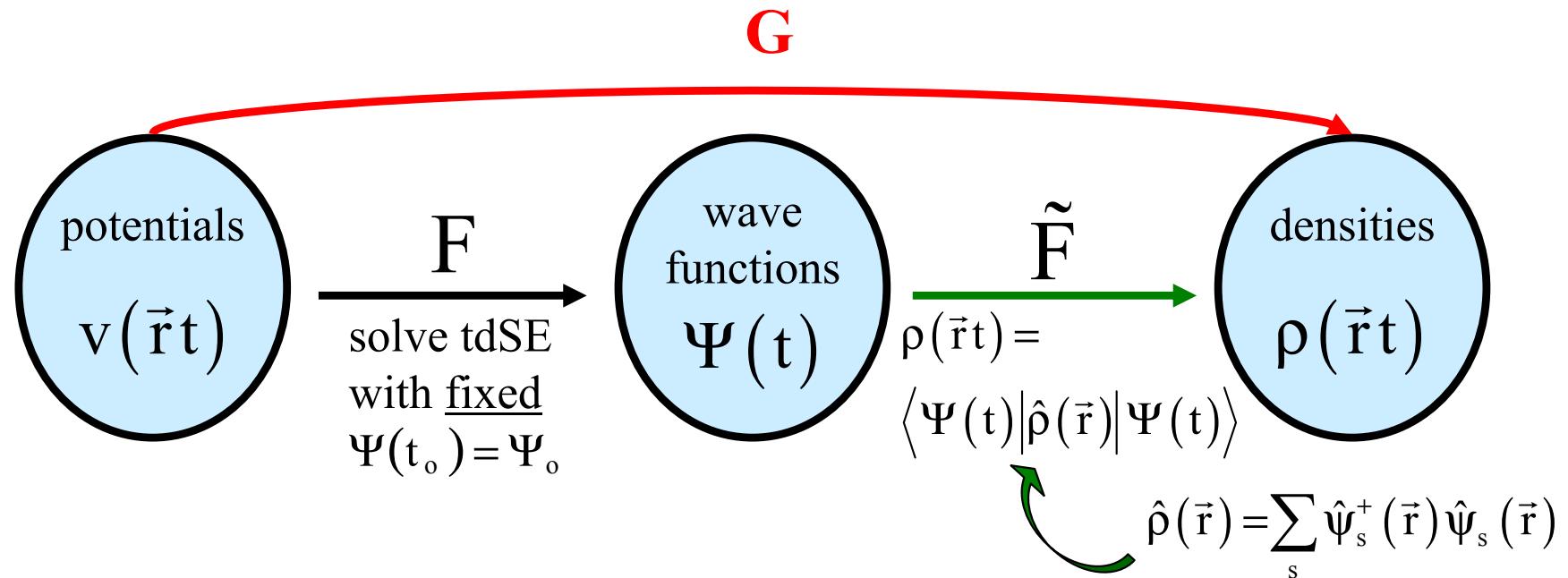
$$G: \quad v(\vec{r}t) \mapsto \rho(\vec{r}t)$$

complete 1 - 1 correspondence <u>not</u> to be expected!

$$i \frac{\partial}{\partial t} \Psi(t) = \left(\hat{T} + \underline{\underline{\hat{V}(t)}} + \hat{W} \right) \Psi(t) \qquad \Psi(t_o) = \Psi_o$$

$$i \frac{\partial}{\partial t} \Psi'(t) = \left(\hat{T} + \underline{\underline{\hat{V}'(t)}} + \hat{W} \right) \Psi'(t) \qquad \Psi'(t_o) = \Psi_o$$

$$\hat{V}'(t) = \hat{V}(t) + C(t) \quad \Leftrightarrow \quad \Psi'(t) = e^{-i\alpha(t)} \Psi(t)$$

"no operator"

with $\boxed{\dot{\alpha}(t) = C(t)}$

$$\Rightarrow \underline{\underline{\rho'(\vec{r}t) = \rho(\vec{r}t)}}$$

i.e. $\left\{ \hat{V}(t) + C(t) \right\} \rightarrow \rho(\vec{r}t)$

If G invertible   up to within time-dependent function C(t)

$$\Rightarrow \Psi = FG^{-1}\rho \qquad \text{fixed up to within time-dependent phase}$$

$$\text{i.e.} \quad \Psi = e^{-i\alpha(t)}\Psi[\rho]$$

For any observable $\hat{O}$

$$\left\langle \Psi \middle| \hat{O} \middle| \Psi \right\rangle = \left\langle \Psi[\rho] \middle| \hat{O} \middle| \Psi[\rho] \right\rangle = O[\rho]$$

is functional of the density

# THEOREM (time-dependent analogue of Hohenberg-Kohn theorem)

The map

$$G: \quad v(\vec{r}t) \mapsto \rho(\vec{r}t)$$

defined for all single-particle potentials $v(\vec{r}t)$ which can be expanded into a Taylor series with respect to the time coordinate around t_o

is invertible up to within an additive merely time-dependent function in the potential.

Proof:
to be shown:

cannot happen
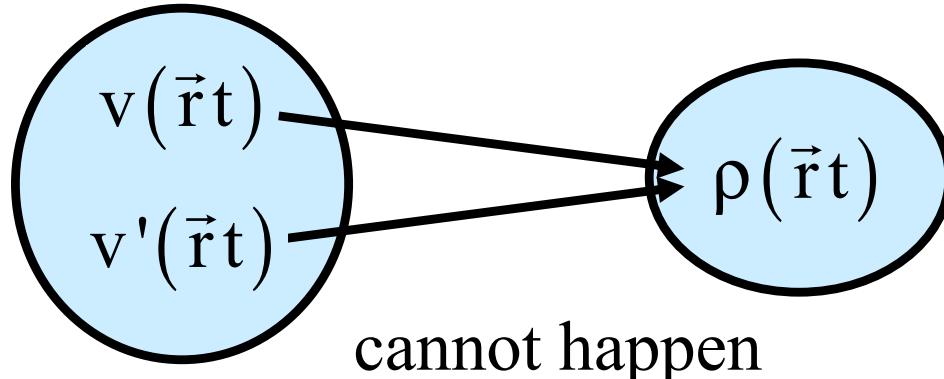
i.e. $\quad \hat{v}(\vec{r}t) \neq \hat{v}'(\vec{r}t) + c(t) \quad \overset{!}{\Rightarrow} \quad \rho(\vec{r}t) \neq \rho'(\vec{r}t)$

potential expandable into Taylor series

$$\exists\, k \geq 0 \quad : \quad \frac{\partial^k}{\partial t^k}\left[v(\vec{r}t) - v'(\vec{r}t)\right]_{t=t_o} \neq \text{constant}$$
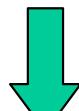
step 1

$$\vec{j}(\vec{r}t) \neq \vec{j}'(\vec{r}t)$$

step 2

$$\rho(\vec{r}t) \neq \rho'(\vec{r}t)$$

Step 1: Current densities

$$\vec{j}(\vec{r}t) = \left\langle \Psi(t) \left| \hat{\vec{j}}(\vec{r}) \right| \Psi(t) \right\rangle$$

$$\text{with} \quad \hat{\vec{j}}(\vec{r}) = -\frac{1}{2i} \sum_s \left(\left[\vec{\nabla} \hat{\psi}_s^+(\vec{r}) \right] \hat{\psi}_s(\vec{r}) - \hat{\psi}_s^+(\vec{r}) \left[\vec{\nabla} \hat{\psi}_s(\vec{r}) \right] \right)$$

Use equation of motion:

$$i\frac{\partial}{\partial t} \left\langle \Psi(t) \left| \hat{O}(t) \right| \Psi(t) \right\rangle = \left\langle \Psi(t) \left| i\frac{\partial}{\partial t}\hat{O}(t) + \left[\hat{O}(t), \hat{H}(t) \right] \right| \Psi(t) \right\rangle$$

$$\Rightarrow \quad i\frac{\partial}{\partial t}\vec{j}(\vec{r}t) = \left\langle \Psi(t) \left| \left[\hat{\vec{j}}(\vec{r}), \hat{H}(t) \right] \right| \Psi(t) \right\rangle$$

$$i\frac{\partial}{\partial t}\vec{j}'(\vec{r}t) = \left\langle \Psi'(t) \left| \left[\hat{\vec{j}}(\vec{r}), \hat{H}'(t) \right] \right| \Psi'(t) \right\rangle$$

note:
$$\vec{j}(\vec{r}\underline{\underline{t_o}}) = \vec{j}'(\vec{r}\underline{\underline{t_o}}) = \left\langle \Psi_o \left| \hat{\vec{j}}(\vec{r}) \right| \Psi_o \right\rangle \equiv \vec{j}_o(\vec{r})$$

$$\rho(\vec{r}\underline{\underline{t_o}}) = \rho'(\vec{r}\underline{\underline{t_o}}) = \left\langle \Psi_o \left| \hat{\rho}(\vec{r}) \right| \Psi_o \right\rangle \equiv \rho_o(\vec{r})$$

$$i\frac{\partial}{\partial t}\left[\vec{j}(\vec{r}t)-\vec{j}'(\vec{r}t)\right]_{t=t_o} = \left\langle \Psi_o \left\| \left[\hat{\vec{j}}(\vec{r}),\hat{H}(t_o)-\hat{H}'(t_o)\right] \right\| \Psi_o \right\rangle$$

$$= \left\langle \Psi_o \left\| \left[\hat{\vec{j}}(\vec{r}),V(t_o)-V'(t_o)\right] \right\| \Psi_o \right\rangle$$

$$= i\rho_o(\vec{r})\vec{\nabla}\left(v(\vec{r}t_o)-v'(\vec{r}t_o)\right)$$

if $\quad \dfrac{\partial^k}{\partial t^k}\left[v(\vec{r}t)-v'(\vec{r}t)\right]_{t=t_o} \neq \text{constant} \quad$ holds for k=0

then $\quad i\dfrac{\partial}{\partial t}\left[\vec{j}(\vec{r}t)-\vec{j}'(\vec{r}t)\right]_{t=t_o} \neq 0$
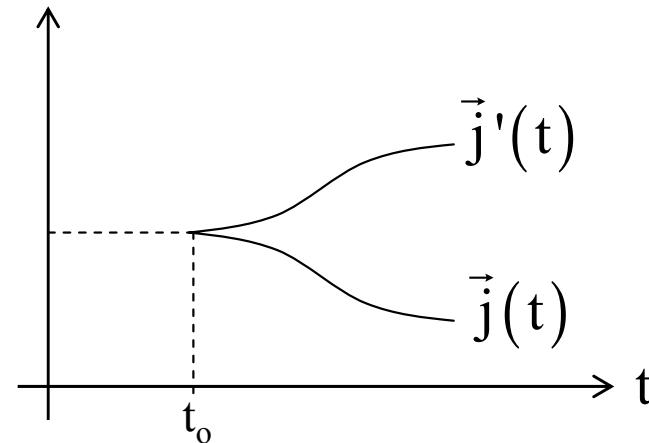
$$\Rightarrow \quad \underline{\underline{\vec{j}(\vec{r}t) \neq \vec{j}'(\vec{r}t)}} \qquad \text{q.e.d.}$$

if $\quad \dfrac{\partial^k}{\partial t^k}\left[v(\vec{r}t)-v'(\vec{r}t)\right]_{t=t_o} \neq$ constant $\qquad$ holds for k>0

$\longrightarrow \quad$ __use equation of motion k+1 times:__

$$\left(i\frac{\partial}{\partial t}\right)^2 \vec{j}(\vec{r}t) = i\frac{\partial}{\partial t}\left\langle \Psi(t)\left|\left[\hat{\vec{j}},\hat{H}(t)\right]\right|\Psi(t)\right\rangle$$

$$= \left\langle \Psi(t)\left|i\frac{\partial}{\partial t}\left[\hat{\vec{j}},\hat{H}(t)\right] + \left[\left[\vec{j},\hat{H}(t)\right],\hat{H}(t)\right]\right|\Psi(t)\right\rangle$$

$$\left(i\frac{\partial}{\partial t}\right)^3 \vec{j}(\vec{r}t) = i\frac{\partial}{\partial t}\left\langle \Psi(t)\left|i\frac{\partial}{\partial t}\left[\vec{j},\hat{H}(t)\right] + \left[\left[\vec{j},H(t)\right],\hat{H}(t)\right]\right|\Psi(t)\right\rangle$$

$$= \cdots$$

$$\left(i\frac{\partial}{\partial t}\right)^{k+1}\left[\vec{j}(\vec{r}t)-\vec{j}\,'(\vec{r}t)\right]_{t=t_o} = i\rho_o(\vec{r})\vec{\nabla}\left(\underbrace{\left(i\frac{\partial}{\partial t}\right)^k\left[v(\vec{r}t)-v'(\vec{r}t)\right]_{t_o}}_{\neq\text{ constant}}\right) \neq 0$$

$$\Rightarrow \quad \underline{\vec{j}(\vec{r}t) \neq \vec{j}\,'(\vec{r}t)} \qquad \text{q.e.d.}$$

Step 2: densities

Use continuity equation:

$$\frac{\partial}{\partial t}\left[\rho(\vec{r}t)-\rho'(\vec{r}t)\right]=-\text{div}\left[\vec{j}(\vec{r}t)-\vec{j}'(\vec{r}t)\right]$$

$$\Rightarrow \quad \frac{\partial^{k+2}}{\partial t^{k+2}}\left[\rho(\vec{r}t)-\rho'(\vec{r}t)\right]_{t=t_o}=-\text{div}\frac{\partial^{k+1}}{\partial t^{k+1}}\left[\vec{j}(\vec{r}t)-\vec{j}'(\vec{r}t)\right]_{t=t_o}$$

$$=-\text{div }\rho_o(\vec{r})\vec{\nabla}\left(\underbrace{\frac{\partial^{k}}{\partial t^{k}}\left[v(\vec{r}t)-v'(\vec{r}t)\right]_{t=t_o}}_{\neq \text{ constant}}\right)$$

remains to be shown:

$$\text{div}\left[\rho_o(\vec{r})\vec{\nabla}u(\vec{r})\right]\neq 0 \qquad \text{if} \qquad u(\vec{r})\neq \text{constant}$$

<u>Proof</u>: by reductio ad absurdum

Assume: $\quad \text{div}\left[\rho_o(\vec{r})\vec{\nabla}u(\vec{r})\right] = 0 \quad$ with $\quad u(\vec{r}) \neq \text{constant}$

$$\int dr^3 \rho_o(\vec{r}) \left(\vec{\nabla}u(\vec{r})\right)^2$$

$$= -\int dr^3 u(\vec{r})\, \underbrace{\text{div}\left[\rho_o(\vec{r})\vec{\nabla}u(\vec{r})\right]}_{0} + \underbrace{\oint \rho_o(\vec{r})u(\vec{r})\vec{\nabla}u(\vec{r})\cdot d\vec{S}}_{0} = 0$$

$$\Rightarrow \quad \rho_o(\vec{r})\left(\vec{\nabla}u(\vec{r})\right)^2 \equiv 0 \quad \longrightarrow \quad \text{contradiction to}$$
$$u(\vec{r}) \neq \text{constant}$$

Simplest possible approximation for $v_{xc}[\rho](\vec{r}t)$

Adiabatic Local Density Approximation (ALDA)

$$v_{xc}^{ALDA}(\vec{r}t) := v_{xc,stat}^{hom}(n)\Big|_{n=\rho(\vec{r}t)}$$

$v_{xc,stat}^{hom}$ = xc potential of static homogeneous e-gas

Approximation with correct asymptotic -1/r behavior: time-dependent optimized effective potential

C. A. Ullrich, U. Gossmann, E.K.U.G., PRL 74, 872 (1995)

LINEAR RESPONSE THEORY

$t = t_0$: **Interacting system in ground state of potential $v_0(r)$ with density $\rho_0(r)$**

$t > t_0$: **Switch on perturbation $v_1(r\ t)$ (with $v_1(r\ t_0) = 0$).**

Density: $\rho(r\ t) = \rho_0(r) + \delta\rho(r\ t)$

Consider functional $\rho[v](r\ t)$ defined by solution of interacting TDSE

Functional Taylor expansion of $\rho[v]$ around v_o:

$$\rho[v]\ (rt) = \rho[v_0 + v_1]\ (rt)$$

$$= \rho[v_0](rt) \qquad\qquad \longrightarrow \rho_o(r)$$

$$+ \int \left.\frac{\delta\rho[v](rt)}{\delta v(r't')}\right|_{v_0} v_1(r't')d^3r'dt' \qquad \longrightarrow \rho_1(rt)$$

$$+ \frac{1}{2}\int\int \left.\frac{\delta^2\rho[v](rt)}{\delta v(r't')\delta v(r''t'')}\right|_{v_0} v_1(r',t')v_1(r'',t'')d^3r'd^3r''dt'dt'' \longrightarrow \rho_2(rt)$$

$$\vdots$$

$\rho_1(r,t) = $ **linear density response of interacting system**

$$\chi(rt, r't') := \left. \frac{\delta\rho[v](rt)}{\delta v(r't')} \right|_{v_0} = $$ **density-density response function of interacting system**

Analogous function $\rho_S[v_S](r\,t)$ for _non_-interacting system

$$\rho_S[v_S](rt) = \rho_S[v_{S,0} + v_{S,1}](rt) = \rho_S[v_{S,0}](rt) + \int \left. \frac{\delta\rho_S[v_S](rt)}{\delta v_S(r't')} \right|_{v_{S,0}} v_{S,1}(r't') \, d^3r'dt' + \cdots$$

$$\chi_S(rt, r't') := \left. \frac{\delta\rho_S[v_S](rt)}{\delta v_S(r't')} \right|_{v_{S,0}} = $$ **density-density response function of _non_-interacting system**

GOAL: Find a way to calculate $\rho_1(r\,t)$ without explicitly evaluating $\chi(r\,t,r't')$ of the <u>interacting</u> system

<u>starting point</u>: **Definition of xc potential**

$$v_{xc}[\rho](rt) := v_S[\rho](rt) - v_{ext}[\rho](rt) - v_H[\rho](rt)$$

<u>Notes</u>:

- v_{xc} is well-defined through non-interacting/ interacting 1-1 mapping.
- $v_S[\rho]$ depends on initial determinant Φ_0.
- $v_{ext}[\rho]$ depends on initial many-body state Ψ_0.

$\Rightarrow$ **In general,** $v_{xc} = v_{xc}[\rho, \Phi_0, \Psi_0]$

only if system is initially in <u>ground-state</u> then, via HK, Φ_0 and Ψ_0 are determined by ρ_0 and v_{xc} depends on ρ alone.

$$\frac{\delta v_{xc}[\rho](rt)}{\delta\rho(r't')}\bigg|_{\rho_0} = \frac{\delta v_S[\rho](rt)}{\delta\rho(r't')}\bigg|_{\rho_0} - \frac{\delta v_{ext}[\rho](rt)}{\delta\rho(r't')}\bigg|_{\rho_0} - \frac{\delta(t-t')}{|r-r'|}$$

$$\uparrow \qquad\qquad \uparrow \qquad\qquad \uparrow \qquad\qquad \uparrow$$

$$f_{xc}(rt,r't') \qquad \chi_S^{-1}(rt,r't') \qquad \chi^{-1}(rt,r't') \qquad W_C(rt,r't')$$

$$\textcolor{red}{\chi_S} \bullet \Big| f_{xc} + W_C = \chi_S^{-1} - \chi^{-1} \Big| \bullet \textcolor{red}{\chi}$$

$$\chi_S\left(f_{xc} + W_C\right)\chi = \chi - \chi_S$$

$$\chi = \chi_S + \chi_S \left(W_C + f_{xc} \right) \chi$$

Act with this operator equation on arbitrary $v_1(r\,t)$ and use $\chi\, v_1 = \rho_1$:

$$\rho_1(rt) = \int d^3r'dt'\chi_S(rt,r't)\left[v_1(rt) + \int d^3r''dt''\left\{ W_C(r't',r''t'') + f_{xc}(r't',r''t'') \right\} \rho_1(r''t'') \right]$$

- Exact integral equation for $\rho_1(r\,t)$, to be solved iteratively

- Need approximation for $f_{xc}(r't',r''t'') = \dfrac{\delta v_{xc}[\rho](r't')}{\delta\rho(r''t'')}\bigg|_{\rho_0}$

  (either for f_{xc} directly or for v_{xc})

Adiabatic approximation

$$v_{xc}^{adiab}[\rho](rt) := v_{xc}^{static\ DFT}[\rho(t)](rt)$$

In the adiabatic approximation, the xc potential $v_{xc}(t)$ at time t only depends on the density $\rho(t)$ at the very same point in time.

e.g. adiabatic LDA: $\quad v_{xc}^{ALDA}(rt) := v_{xc}^{LDA}(\rho(rt)) = -\alpha\,\rho(rt)^{1/3} + \cdots$

$$\Rightarrow f_{xc}^{ALDA}(rt, r't') = \left.\frac{\delta v_{xc}^{ALDA}(rt)}{\delta\rho(r't')}\right|_{\rho_0} = \delta(r-r')\delta(t-t')\left.\frac{\partial v_{xc}^{ALDA}}{\partial\rho(r)}\right|_{\rho_0(r)}$$

$$= \delta(r-r')\delta(t-t')\left.\frac{\partial^2 e_{xc}^{hom}}{\partial n^2}\right|_{\rho_0(r)}$$

Total photoabsorption cross section of the Xe atom versus photon energy in the vicinity of the 4d threshold.
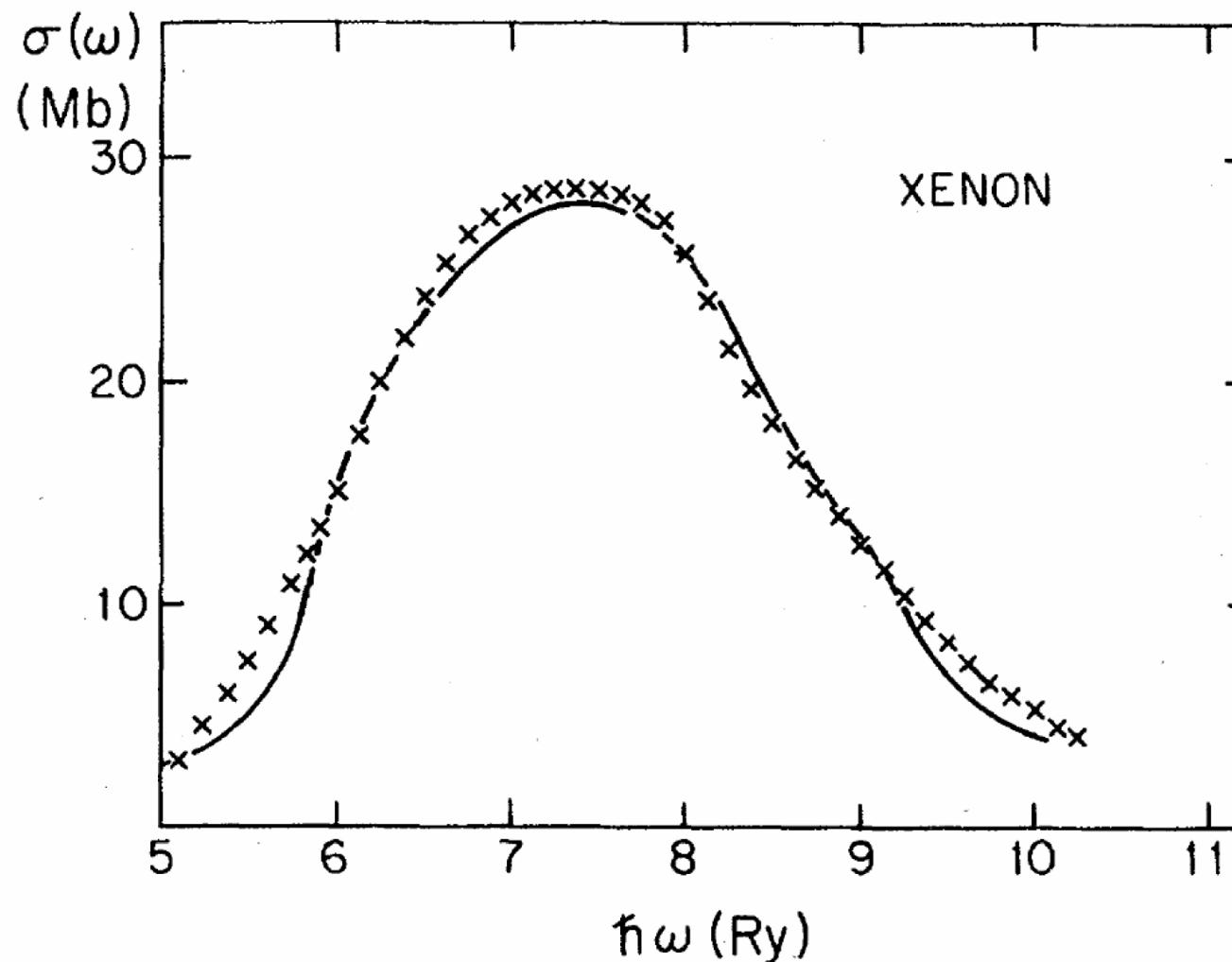
Solid line: self-consistent time-dependent KS calculation [A. Zangwill and P. Soven, Phys. Rev. A **21**, 1561 (1980)]; crosses: experimental data [R. Haensel, G. Keitel, P. Schreiber, and C. Kunz, Phys. Rev. **188**, 1375 (1969)].

Photo-absorption in weak lasers
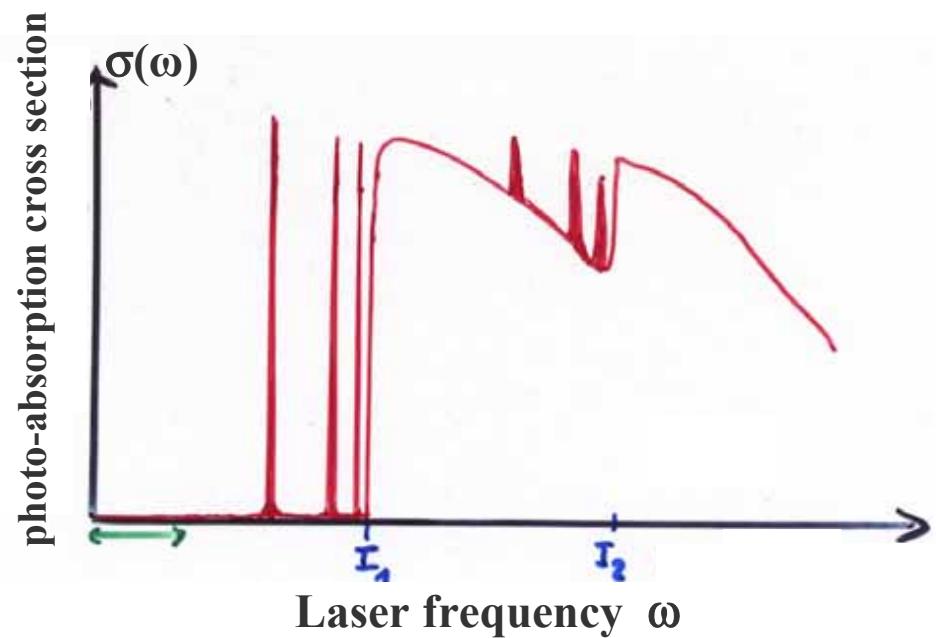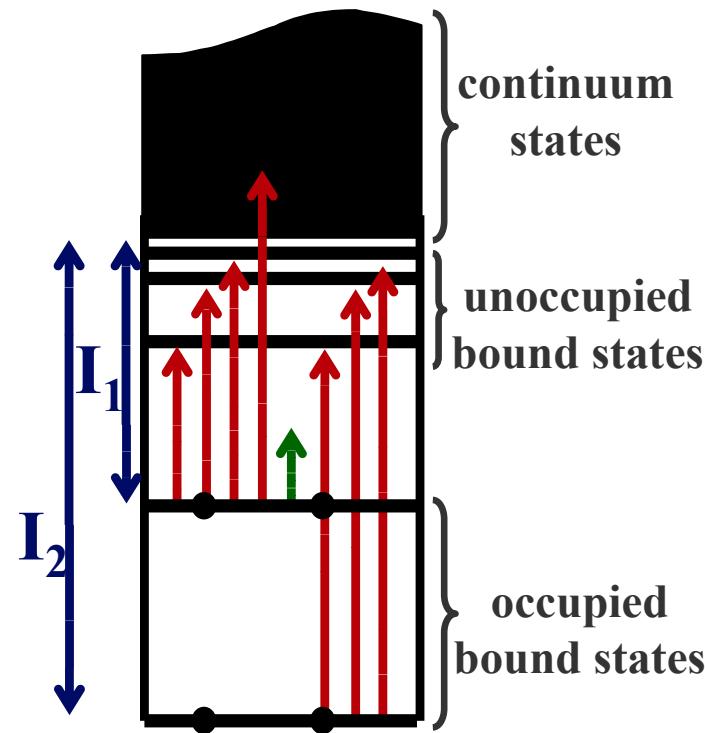
No absorption if ω < lowest excitation energy

Standard linear response formalism

H(t$_0$) = full static Hamiltonian at t$_0$

$$H(t_0)|m\rangle = E_m|m\rangle \quad \leftarrow \text{ exact many-body eigenfunctions}$$
$$\text{and energies of system}$$

full response function

$$\chi(r,r';\omega) = \lim_{\eta \to 0^+} \sum_m \left(\frac{\langle 0|\hat{\rho}(r)|m\rangle\langle m|\hat{\rho}(r)|0\rangle}{\omega - (E_m - E_0) + i\eta} - \frac{\langle 0|\hat{\rho}(r')|m\rangle\ \langle m|\hat{\rho}(r')|0\rangle}{\omega + (E_m - E_0) + i\eta} \right)$$

$\Rightarrow$ **The exact linear density response**

$$\rho_1(\omega) = \hat{\chi}(\omega) v_1(\omega)$$

has poles at the exact excitation energies $\Omega = E_m - E_0$

Discrete excitation energies from TDDFT

exact representation of linear density response:

$$\rho_1(\omega) = \hat{\chi}_S(\omega)\left(v_1(\omega) + \hat{W}_C \rho_1(\omega) + \hat{f}_{xc}(\omega)\rho_1(\omega)\right)$$

"^" denotes integral operators, i.e. $\hat{f}_{xc}\rho_1 \equiv \int f_{xc}(\vec{r},\vec{r}')\rho_1(\vec{r}')d^3r'$

where $\hat{\chi}_S(\vec{r},\vec{r}';\omega) = \sum\limits_{j,k} \dfrac{M_{jk}(\vec{r},\vec{r}')}{\omega - (\varepsilon_j - \varepsilon_k) + i\eta}$

with $M_{jk}(\vec{r},\vec{r}') = (f_k - f_j)\varphi_j(\vec{r})\varphi_j^*(\vec{r}')\varphi_k(\vec{r}')\varphi_k^*(\vec{r})$

$$f_m = \begin{cases} 1 \text{ if } \varphi_m \text{ is occupied in KS ground state} \\ 0 \text{ if } \varphi_m \text{ is unoccupied in KS ground state} \end{cases}$$

$\varepsilon_j - \varepsilon_k$ KS excitation energy

$$\left(\hat{1} - \hat{\chi}_S(\omega) \left[\hat{W}_C + \hat{f}_{xc}(\omega) \right] \right) \rho_1(\omega) = \hat{\chi}_S(\omega) v_1(\omega)$$

$\rho_1(\omega) \to \infty$ for $\omega \to \Omega$ (exact excitation energy) but right-hand side remains finite for $\omega \to \Omega$

hence $\left(\hat{1} - \hat{\chi}_S(\omega) \left[\hat{W}_C + \hat{f}_{xc}(\omega) \right] \right) \xi(\omega) = \lambda(\omega) \xi(\omega)$

$\lambda(\omega) \to 0$ for $\omega \to \Omega$

This condition rigorously determines the exact excitation energies, i.e.,

$$\left(\hat{1} - \hat{\chi}_S(\Omega) \left[\hat{W}_C + \hat{f}_{xc}(\Omega) \right] \right) \xi(\Omega) = 0$$

This leads to the (non-linear) eigenvalue equation

(See T. Grabo, M. Petersilka, E. K. U. G., J. Mol. Struc. (Theochem) <u>501</u>, 353 (2000))

$$\sum_{q'} \left(A_{qq'}(\Omega) + \omega_q \delta_{qq'} \right) \beta_{q'} = \Omega \beta_q$$

where

$$A_{qq'} = \alpha_{q'} \int d^3r \int d^3r' \Phi_q(r) \left(\frac{1}{|r - r'|} + f_{xc}(r, r', \Omega) \right) \Phi_{q'}(r')$$

$$q = (j, a) \text{ double index} \qquad \alpha_q = f_a - f_j$$

$$\Phi_q(r) = \varphi_a^*(r) \varphi_j(r) \qquad \omega_q = \varepsilon_a - \varepsilon_j$$

Atom	Experimental Excitation Energies $^1S \to {}^1P$ (in Ry)	KS energy differences $\Delta\epsilon_{KS}$ (Ry)	$\Delta\epsilon_{KS} + K$
Be	0.388	0.259	0.391
Mg	0.319	0.234	0.327
Ca	0.216	0.157	0.234
Zn	0.426	0.315	0.423
Sr	0.198	0.141	0.210
Cd	0.398	0.269	0.391

from: M. Petersilka, U. J. Gossmann, E.K.U.G., PRL 76, 1212 (1996)

$$\Delta E = \underbrace{\Delta\epsilon_{KS}}_{\epsilon_j - \epsilon_k} + K$$

$$K = \int d^3r \int d^3r' \, \varphi_j(r)\varphi_j^*(r')\varphi_k(r')\varphi_k^*(r)\left(\frac{1}{|r-r'|} + f_{xc}(r,r')\right)$$

Excitation energies of CO molecule

State		Ω_{expt}	KS-transition	$\Delta\in_{KS}$	$\Delta\in_{KS} + K$
A	$^1\Pi$	0.3127	$5\Sigma \rightarrow 2\Pi$	0.2523	0.3267
a	$^3\Pi$	0.2323			0.2238
I	$^1\Sigma^-$	0.3631	$1\Pi \rightarrow 2\Pi$	0.3626	0.3626
D	$^1\Delta$	0.3759			0.3812
a'	$^3\Sigma^+$	0.3127			0.3181
e	$^3\Sigma^-$	0.3631			0.3626
d	$^3\Delta$	0.3440			0.3404

approximations made: v_{xc}^{LDA} and f_{xc}^{ALDA}

Quantum defects in Helium $E_n = -\dfrac{1}{2(n-\mu_n)^2}$ [a.u.]

3P Series

Legend:
- Bare KS
- exact
- x-ALDA
- xc-ALDA (VWN)
- x-TDOEP
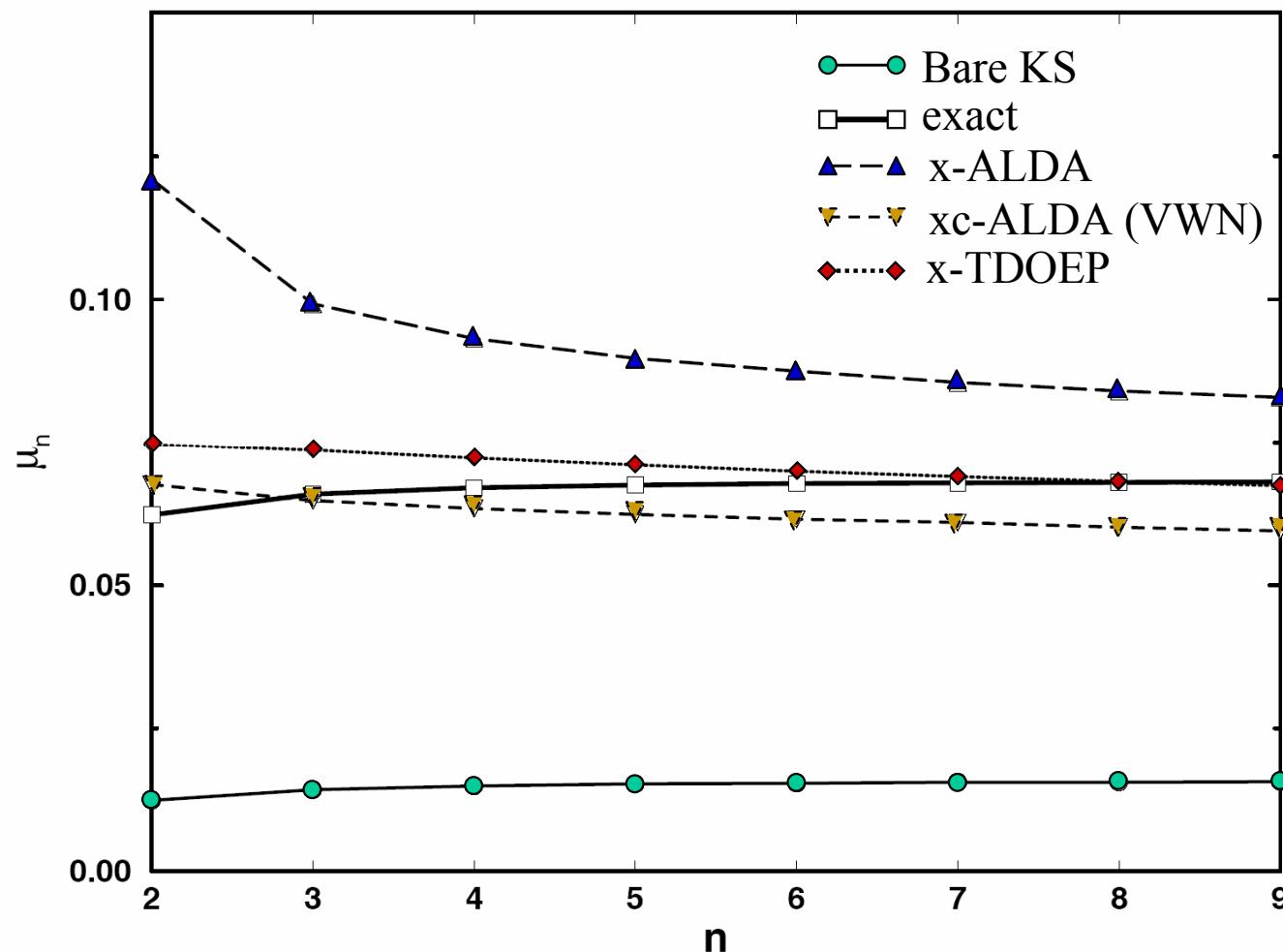
M. Petersilka, U.J. Gossmann and E.K.U.G., in: Electronic Density Functional Theory: Recent Progress and New Directions, J.F. Dobson, G. Vignale, M.P. Das, ed(s), (Plenum, New York, 1998), p 177 - 197.
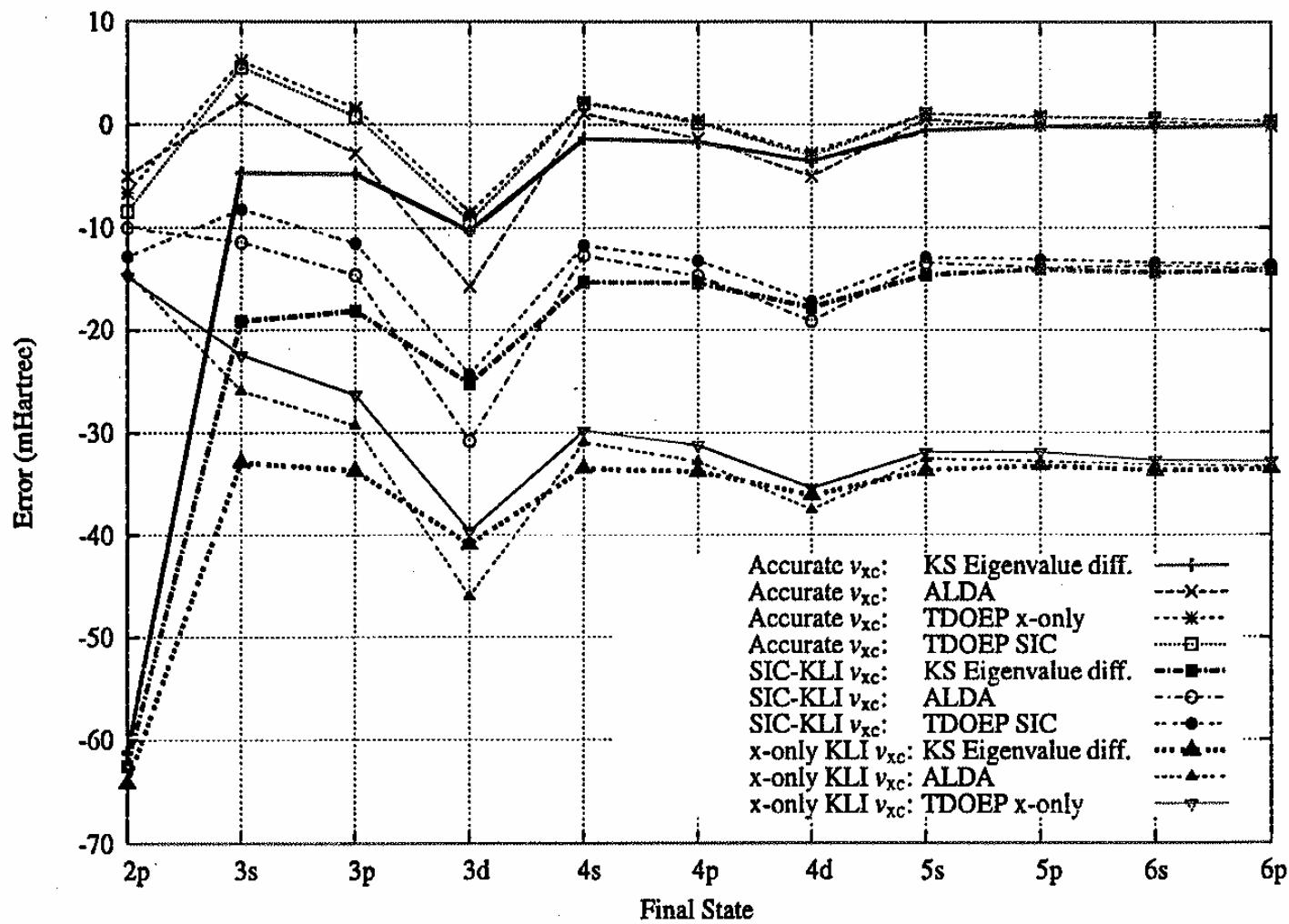
Figure 3.3: Errors of singlet excitation energies from the ground state of Be, calculated from the accurate, the OEP-SIC and x-only KLI exchange correlation potential and with different approximations for the exchange-correlation kernel (see text). The errors are given in mHartrees. To guide the eye, the errors of the discrete excitation energies were connected with lines.

M. Petersilka, E.K.U.G., K. Burke, Int. J. Quantum Chem. **80**, 534 (2000)

Failures of ALDA in the linear response regime

- **H_2 dissociation is incorrect:**

$$E\left(^1\Sigma_u^+\right) - E\left(^1\Sigma_g^+\right) \xrightarrow[R\to\infty]{} 0 \text{ (in ALDA)}$$

(see: Gritsenko, van Gisbergen, Görling, Baerends, J. Chem. Phys. 113, 8478 (2000))

- **response of long chains strongly overestimated**

(see: Champagne et al., J. Chem. Phys. 109, 10489 (1998) and 110, 11664 (1999))

- **in periodic solids, $f_{xc}^{ALDA}\left(q,\omega,\rho\right) = c\left(\rho\right)$ whereas,**

for insulators, $f_{xc}^{exact} \xrightarrow[q\to 0]{} 1/q^2$ divergent.

- **charge-transfer excitations not properly described**

(see: Dreuw et al., J. Chem. Phys. 119, 2943 (2003))

How good is ALDA for solids?

optical absorption (q=0)
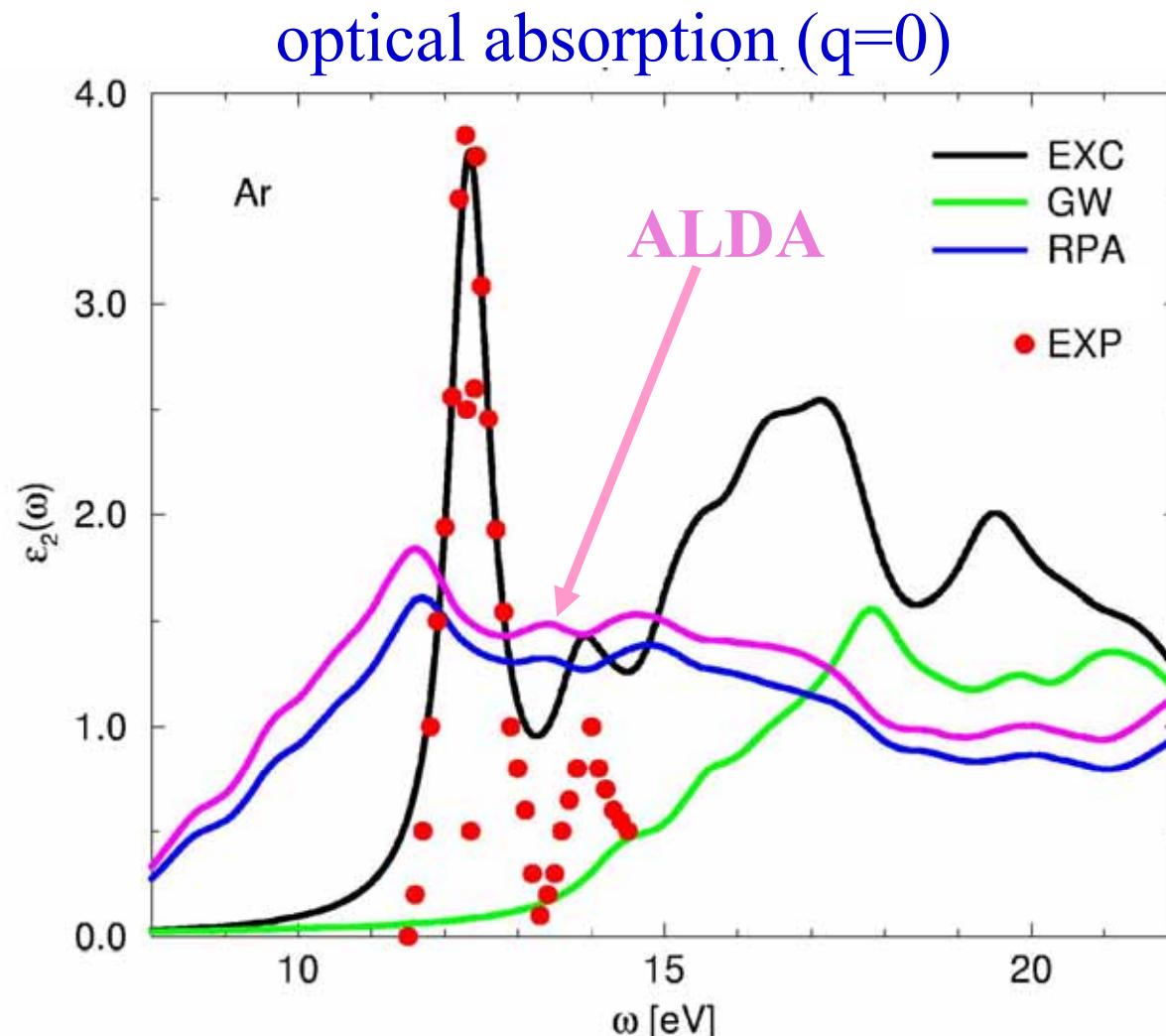
Solid Argon

OBSERVATION:
In the long-wavelength-limit (q = 0), relevent for optical absorption, ALDA is not reliable. In particular, excitonic lines are completely missed. Results are very close to RPA.

EXPLANATION:
In the TDDFT response equation, the bare Coulomb interaction and the xc kernel only appear as sum ($W_C + f_{xc}$). For $q \to 0$, W_C diverges like $1/q^2$, while f_{xc} in ALDA goes to a constant. Hence results are close to $f_{xc} = 0$ (RPA) in the $q \to 0$ limit.

CONCLUSION:
Approximations for f_{xc} are needed which, for $q \to 0$, correctly diverge like $1/q^2$. Such approximations can be derived from many-body perturbation theory (see, e.g., L. Reining, V. Olevano, A. Rubio, G. Onida, PRL 88, 066404 (2002)).

X-ray absorption spectroscopy of 3d metals

d band

E_F

Ω_2 Ω_1

2p$_{3/2}$ (4-fold degenerate $M = -\frac{3}{2}, -\frac{1}{2}, +\frac{1}{2}, +\frac{3}{2}$)

2p$_{1/2}$ (2-fold degenerate $M = -\frac{1}{2}, +\frac{1}{2}$)
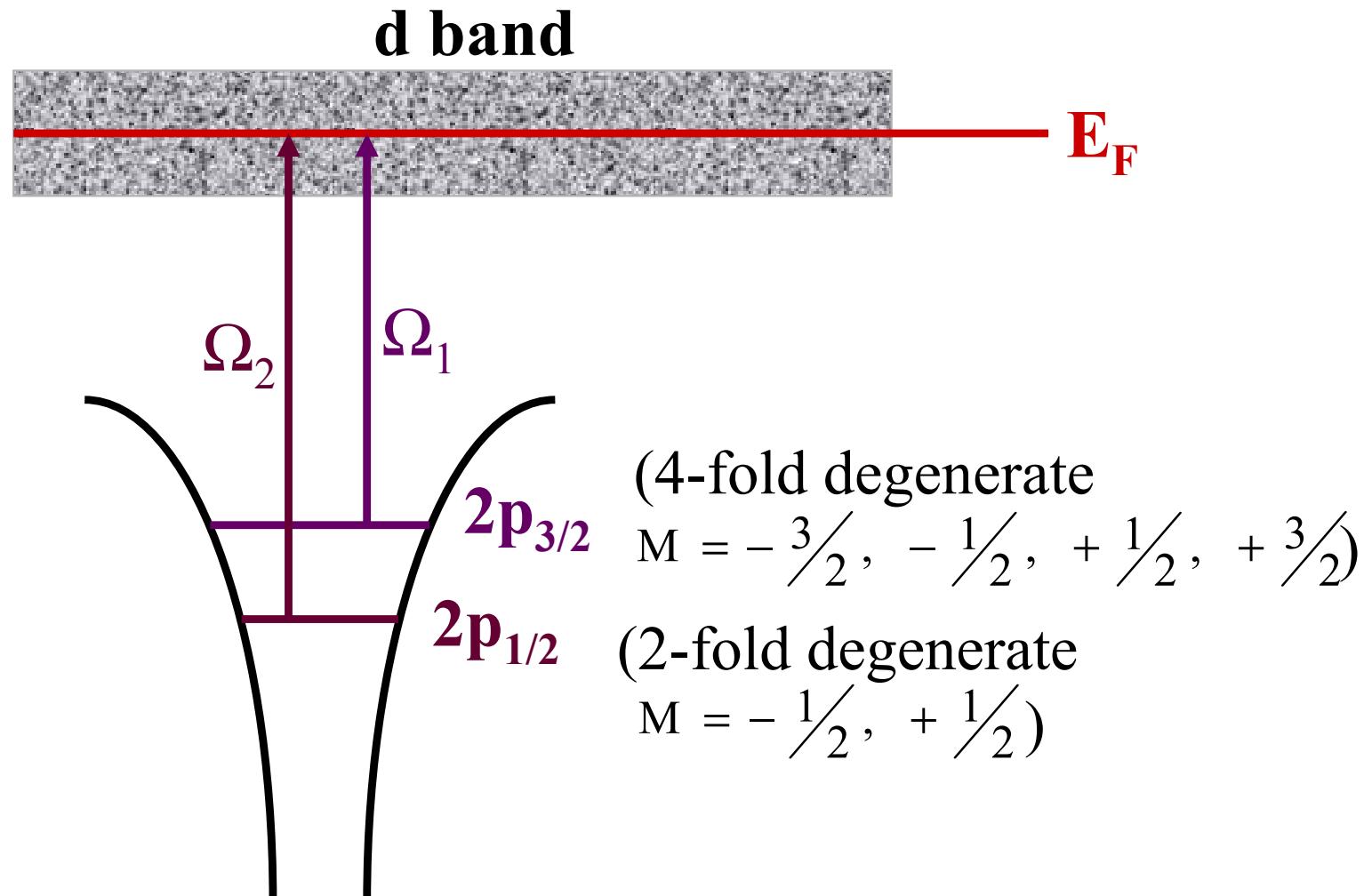
non-interacting (KS) response: $\chi_s \approx \dfrac{c_1}{\omega - \omega_1} + \dfrac{c_2}{\omega - \omega_2}$

Two-pole approximation

$$\begin{pmatrix} \omega_1 + M_{11} & M_{12} \\ M_{21} & \omega_2 + M_{22} \end{pmatrix} \beta = \Omega \beta$$

yields excellent agreement with experiment.

Furthermore, from knowledge of the KS orbitals and the KS excitation energies ω_1, ω_2 and the experimental excitation energies Ω_1, Ω_2 and their branching ratio, one can deduce experimental values for the matrix elements M_{ij}, i.e. one can "measure" f_{xc}.

A. Scherz, E.K.U.G., H. Appel, C. Sorg, K. Baberschke, H. Wende, K. Burke, PRL 95, 253006 (2005)
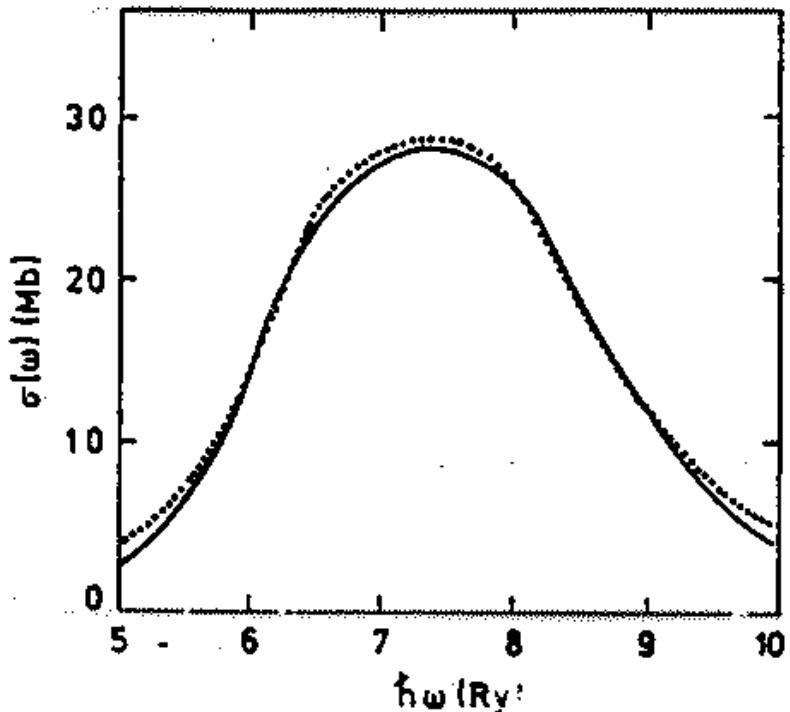
Total photo-absorption cross section for Xenon above the 4d-threshold (Zangwill & Soven 1980)

Approximation used for f_{xc}:
Adiabatic LDA, i.e. $v_{xc}^{ALDA}(\vec{r}, t) = \dfrac{d}{d\rho}\left(\rho \varepsilon_{xc}^{hom}(\rho)\right)\Bigg|_{\rho = \rho(\vec{r}, t)}$

$\Rightarrow$

$$f_{xc}^{ALDA}(\vec{r}, \vec{r}'; \omega) = \delta(\vec{r} - \vec{r}')\dfrac{d^2}{d\rho^2}\left(\rho \varepsilon_{xc}^{hom}(\rho)\right)\Bigg|_{\rho = \rho_0(\vec{r})} = \delta(\vec{r} - \vec{r}') f_{xc}^{hom}(\rho, q = 0, \omega = 0)\Bigg|_{\rho = \rho_0(\vec{r})}$$

no frequency dependence!

Properties of exact $f_{xc}^{hom}(q, \omega)$

1. $f_{xc}^{hom}(q = 0, \omega = 0) = \dfrac{d^2}{d\rho^2}\big(\rho \cdot \varepsilon_{xc}(\rho)\big) \equiv f_0(\rho)$ **(compressibility sum rule)**

2. $f_{xc}^{hom}(q = 0, \omega = \infty) = -\dfrac{4}{5}\rho^{2/3}\dfrac{d}{d\rho}\left(\dfrac{\varepsilon_{xc}(\rho)}{\rho^{2/3}}\right) + 6\rho^{1/3}\dfrac{d}{d\rho}\left(\dfrac{\varepsilon_{xc}(\rho)}{\rho^{1/3}}\right)$

$\equiv f_\infty(\rho)$ **(3rd frequency-moment sum rule)**

3. $f_0(\rho) < f_\infty(\rho) < 0$ **for all ρ** **(according to best approximations known for $\varepsilon_{xc}(\rho)$)**

4. $\operatorname{Re} f_{xc}^{hom}(q, \omega) = \operatorname{Re} f_{xc}^{hom}(q, -\omega)$

$\operatorname{Im} f_{xc}^{hom}(q, \omega) = -\operatorname{Im} f_{xc}^{hom}(q, -\omega)$

$$\boxed{\textbf{Properties of exact } f_{xc}^{hom}(q, \omega)}$$

5. Kramers - Kronig relations

$$\text{Re } f_{xc}^{hom}(q, \omega) - f_{xc}^{hom}(q, \infty) = P \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \frac{\text{Im } f_{xc}^{hom}(q, \omega')}{\omega - \omega'}$$

$$\text{Im } f_{xc}^{hom}(q, \omega) = -P \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \frac{\text{Re } f_{xc}(q, \omega') - f_{xc}(q, \infty)}{\omega - \omega'}$$

6. $\quad \text{Im } f_{xc}^{hom}(q = 0, \omega) \xrightarrow[\omega \to \infty]{} -\dfrac{c}{\omega^{3/2}} \quad, \quad c = \dfrac{23\pi}{15}$

    (follows from perturbation expansion of irreducible polarization to 2nd order in e²)

7. $\quad \text{Re } f_{xc}^{hom}(q = 0, \omega) \xrightarrow[\omega \to \infty]{} f_{\infty} + \dfrac{c}{\omega^{3/2}} \quad$ **(from 5 and 6)**

Parametrization for long wavelength limit

$$\mathrm{Im}\, f_{xc}(q=0,\omega) = \frac{\alpha\omega}{(1+\beta\omega)^{5/4}}$$

with $\quad \alpha(\rho) = -A\left(f_\infty(\rho) - f_0(\rho)\right)^{5/3} \qquad \beta(\rho) = B\left(f_\infty(\rho) - f_0(\rho)\right)^{4/3}$

$\qquad$ A, B > 0 independent of ρ

$\overset{KK}{\Longrightarrow}$

$$\mathbf{Re}\, \mathbf{f}_{xc}(q=0,\omega) = f_\infty + \frac{2\sqrt{2}\,\alpha}{\pi\sqrt{\beta}\,r^2}\cdot\left[2E\left(\frac{1}{\sqrt{2}}\right) - \frac{1+r}{2}\prod\left(\frac{1-r}{2},\frac{1}{\sqrt{2}}\right) - \frac{1-r}{2}\prod\left(\frac{1+r}{2},\frac{1}{\sqrt{2}}\right)\right]$$

with $\quad r := \sqrt{1+\beta\omega^2}$, E, Π = elliptic integrals of 2nd, 3rd kind
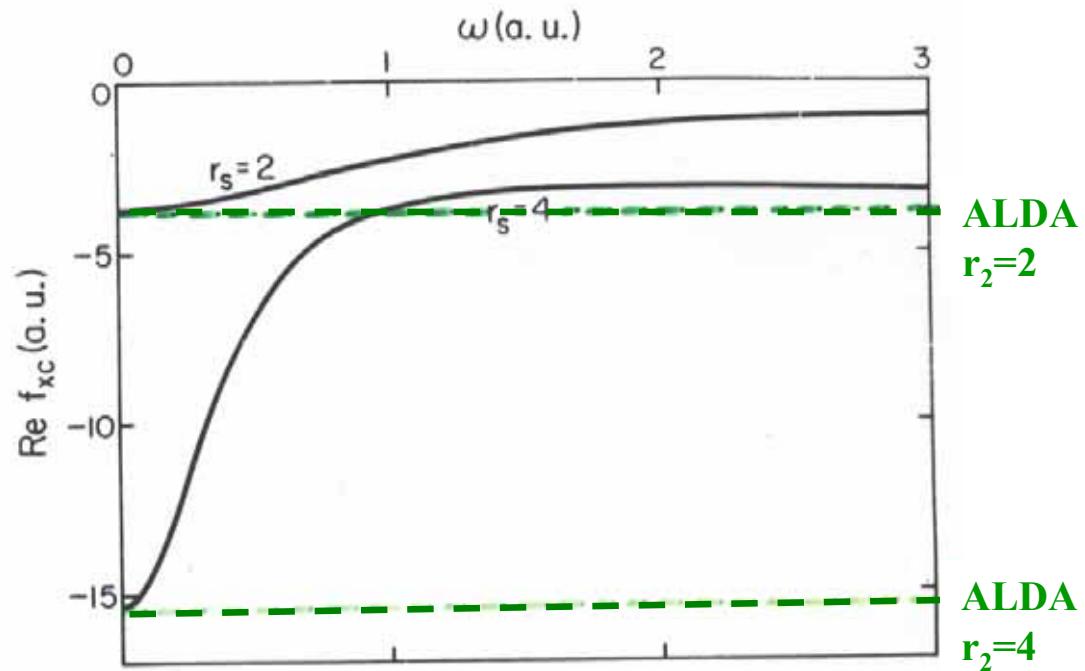
ALDA
$r_2=2$

ALDA
$r_2=4$

ALDA
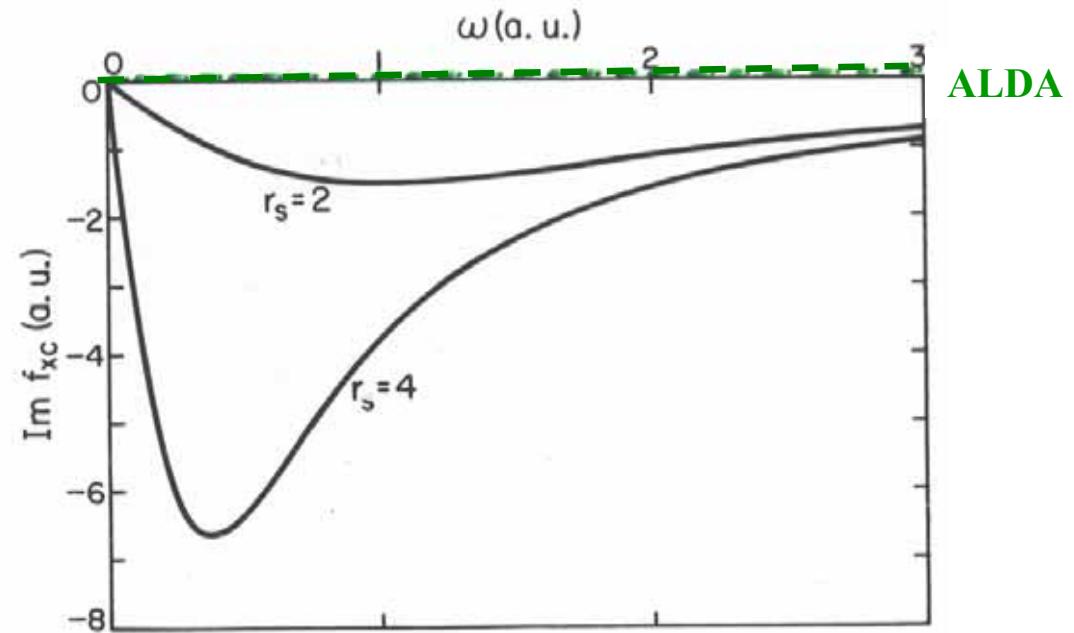
E.K.U.G., W. Kohn, PRL 55, 2850 (1985)
N. Iwamoto, E.K.U.G., Phys. Rev. B 35, 3003 (1987)